

# Energy Conservation in a Pulley + Mass System

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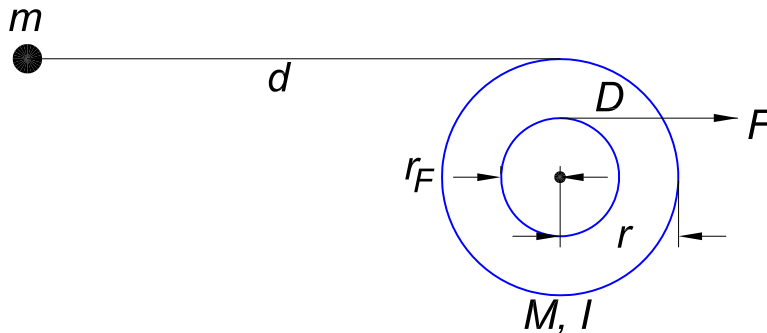
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## 1 Problem

Discuss the motion in the frictionless system of a pulley of mass  $M$ , and moment of inertia  $I$  about its axis, when constant force  $F$  is applied via a string at radius  $r_F$  from the axis of pulley, while a mass  $m$  is also connected to the pulley via a string at radius  $r$ . Verify that the kinetic energy of the system, after starting from rest, equals the work done by force  $F$ .

The translational motion is entirely in the direction of the strings, as sketched below.



*This problem was suggested by Mario Carvajal.*

## 2 Solution

### 2.1 Massless Pulley with Fixed Axis

We first consider a simpler case, of a massless pulley with fixed axis.

The massless pulley has zero moment of inertia  $I$ , so the torque equation for the pulley is

$$\tau = I\alpha = 0 = r_F F - r f, \quad f = \frac{r_F}{r} F. \quad (1)$$

where  $f$  is the (constant) tension in the string that connects mass  $m$  to the pulley. Note that  $f = F$  only if  $r = r_F$ .

As the pulley rotates through angle  $\theta$ , the end of the string supporting force  $F$  advances by distance  $\Delta x_F = r_F \theta$ , while mass  $m$  is pulled towards the pulley by distance  $\Delta x_m = r \theta$ . That is,

$$\theta = \frac{\Delta x_F}{r_F} = \frac{\Delta x_m}{r}, \quad \Delta x_m = \frac{r}{r_F} \Delta x_F. \quad (2)$$

Supposing force  $F$  is first applied at time  $t = 0$ , when mass  $m$  is at rest, the latter accelerates according to

$$a_m = \frac{f}{m}, \quad v_m = \frac{f}{m}t, \quad \Delta x_m = \frac{f}{2m}t^2. \quad (3)$$

For completeness, we note that the end of the string that supports force  $F$  moves according to

$$\Delta x_F = \frac{r_F}{r}\Delta x_m = \frac{r_F}{r}\frac{f}{2m}t^2 = \frac{r_F^2}{r^2}\frac{F}{2m}t^2, \quad v_F = \frac{r_F^2}{r^2}\frac{F}{m}t, \quad a_F = \frac{r_F^2}{r^2}\frac{F}{m}. \quad (4)$$

The kinetic energy of mass  $m$  at time  $t$  is, using eqs. (1)-(3),

$$\text{KE}_m = \frac{m v_m^2}{2} = \frac{m}{2}\frac{f^2}{m^2}t^2 = f\frac{f}{2m}t^2 = f\Delta x_m = \frac{r_F}{r}F\Delta x_m = F\Delta x_F = W_F, \quad (5)$$

such that the work  $W_F = F\Delta x_F$  done by force  $F$  equals the (change of) kinetic energy of mass  $m$ .

## 2.2 Massive Pulley with Fixed Axis

In this case, the torque equation for the pulley becomes

$$\tau = I\alpha = I\ddot{\theta} = r_FF - rf, \quad \omega = \dot{\theta} = \frac{r_FF - rf}{I}t, \quad \theta = \frac{r_FF - rf}{2I}t^2. \quad (6)$$

Equations (2)-(3) remain the same, while the kinetic energy of the system includes the rotational energy of the pulley,

$$\begin{aligned} \text{KE} &= \frac{m v_m^2}{2} + \frac{I\omega^2}{2} = \frac{m}{2}\frac{f^2}{m^2}t^2 + \frac{(r_FF - rf)^2}{2I}t^2 = f\frac{f}{2m}t^2 + (r_FF - rf)\theta \\ &= f\Delta x_m + F\Delta x_F - f\Delta x_m = F\Delta x_F = W_F. \end{aligned} \quad (7)$$

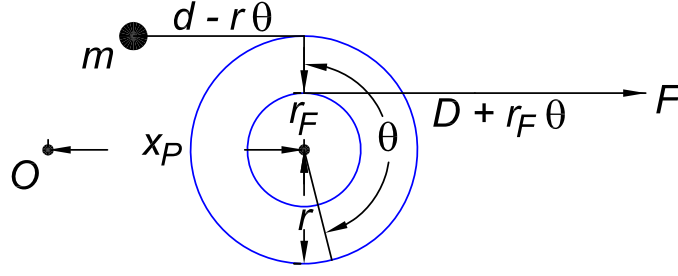
Again, the work done by force  $F$  appears as the kinetic energy of the system.

## 2.3 Massive Pulley whose Axis is not Fixed

We take the center of the pulley to be at the origin at time  $t = 0$ , when force  $F$  is first applied, in the  $-x$  direction, and the end of the string is at distance  $D$  from the pulley. Also at  $t = 0$ , the length of the string between the pulley and mass  $m$  is  $d$ , such that the initial  $x$ -coordinate of the mass is  $x_{m,0} = -d$ .

At a later time  $t$ , the center of the pulley is at  $x_P$  and the pulley has rotated by angle  $\theta > 0$ . The string by which force  $F$  is applied now has unwrapped by length  $r_F\theta$ , and its end has coordinate  $x_F = x_P + D + r_F\theta$ . Similarly, the coordinate of mass  $m$  at time  $t$  is  $x_m = x_P - d + r\theta$ , since length  $r\theta$  of its string is now wrapped around the pulley. That is,

$$\Delta x_F = x_P + r_F\theta, \quad \Delta x_m = x_P + r\theta. \quad (8)$$



Denoting the tension in the string connected to mass  $m$  by  $f > 0$ , the equations of motion for the mass and pulley are

$$m a = m \ddot{x}_m = m(\ddot{x}_P + r \ddot{\theta}) = f, \quad (9)$$

$$M a_P = M \ddot{x}_P = F - f, \quad (10)$$

$$I \alpha = I \ddot{\theta} = \tau = r_F F - r f. \quad (11)$$

These are three equations for the three unknowns  $x_P$ ,  $\theta$  and  $f$ .

Combining eqs. (9) and (10), we find

$$f = \frac{m}{M} (F - f) + m r \ddot{\theta}, \quad f = \frac{m}{M + m} F + \frac{M m}{M + m} r \ddot{\theta}, \quad (12)$$

and using this in eq. (11) leads to

$$I \ddot{\theta} = r_F F - \frac{m}{M + m} r F - \frac{M m}{M + m} r^2 \ddot{\theta}, \quad I' \ddot{\theta} = R F, \quad (13)$$

where

$$I' = I + \frac{M m r^2}{M + m}, \quad R = r_F - \frac{m}{M + m} r. \quad (14)$$

For  $r/r_F > (M + m)/m$ , the quantity  $R$  is negative, and the motion would not be as in the figure above; instead the pulley would rotate “backwards”, and its center would move away from the  $x$ -axis.

We restrict further discussion to the case that  $R > 0$ .

From eq. (13), we have that

$$\omega \equiv \dot{\theta} = \frac{R F}{I'} t, \quad \theta = \frac{R F}{2 I'} t^2. \quad (15)$$

The tension in the string connected to mass  $m$  follows from eqs. (12)-(13) as

$$f = \frac{m}{M + m} \left( 1 + \frac{M r R}{I'} \right) F = \frac{m}{M + m} \frac{I + M r r_F}{I + \frac{M m r^2}{M + m}} F, \quad (16)$$

which is a constant force, with  $f \leq F$  for  $R > 0$ .

The motion of mass  $m$  follows from eq. (9) as

$$v_m \equiv \dot{x}_m = \frac{f}{m} t, \quad \Delta x_m = \frac{f}{2 m} t^2, \quad (17)$$

and the motion of the center of the pulley follows from eq. (10) as

$$v_P \equiv \dot{x}_P = \frac{F-f}{M}t, \quad x_P = \frac{F-f}{2M}t^2. \quad (18)$$

The kinetic energy of the system is

$$\begin{aligned} \text{KE} &= \frac{m v_m^2}{2} + \frac{M v_P^2}{2} + \frac{I \omega^2}{2} = \frac{f^2 t^2}{2m} + \frac{(F-f)^2 t^2}{2M} + I \frac{R^2 F^2 t^2}{2I'^2} \\ &= f \Delta x_m + (F-f)x_P + F \frac{IR}{I'} \theta = f(x_p + r\theta) + (F-f)x_P + F \frac{IR}{I'} \theta \\ &= Fx_P + f r \theta + F \frac{IR}{I'} \theta = Fx_P + F \frac{mr}{M+m} \left(1 + \frac{MrR}{I'}\right) \theta + F \frac{IR}{I'} \theta \\ &= Fx_P + F(r_F - R) \left(1 + \frac{MrR}{I'}\right) \theta + F \frac{IR}{I'} \theta \\ &= F(x_P + r_F \theta) + FR\theta \left[ r_F \frac{Mr}{I'} - \left(1 + \frac{MrR}{I'}\right) + \frac{I}{I'} \right] \\ &= F \Delta x_F + FR\theta \left[ (r_F - R) \frac{Mr}{I'} + \frac{I}{I'} - 1 \right] \\ &= W_F + FR\theta \left[ \frac{m}{M+m} \frac{Mr^2}{I'} + \frac{I}{I'} - 1 \right] = W_F. \end{aligned} \quad (19)$$

Once again, the work done by force  $F$  appears as the kinetic energy of the system.